A Limit on the Speed of Quantum Computation in Determining Parity $^{\circ}$

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Consider a function f which is defined on the integers from 1 to N and takes the values -1 and +1. The parity of f is the product over all x from 1 to N of f(x). With no further information about f, to classically determine the parity of f requires N calls of the function f. We show that any quantum algorithm capable of determining the parity of f contains at least N/2 applications of the unitary operator which evaluates f. Thus for this problem, quantum computers cannot outperform classical computers.

I. INTRODUCTION

If a quantum computer is ever built, it could be used to solve certain problems in less time than a classical computer. Simon found a problem that can be solved exponentially faster by a quantum computer than by the provably best classical algorithm [1]. The Shor algorithm for factoring on a quantum computer gives an exponential speedup over the best known classical algorithm [2]. The Grover algorithm gives a speedup for the following problem [3]. Suppose you are given a function f(x) with x an integer and $1 \le x \le N$. Furthermore you know that f is either identically equal to 1 or it is 1 for N-1 of the x's and equal to -1 at one unknown value of x. The task is to determine which type of f you have. Without any additional information about f, classically this takes of order N calls of f whereas the quantum algorithm runs in time of order \sqrt{N} . In fact this \sqrt{N} speedup can be shown to be optimal [4].

It is of great interest to understand the circumstances under which quantum speedup is possible. Recently Ozhigov has shown that there is a situation where a quantum computer cannot outperform a classical computer [5]. Consider a function g(t), defined on the integers from 1 to L, which takes integer values from 1 to L. We wish to find the M^{th} iterate of some input, say 1, that is, $g^{[M]}(1)$. (Here $g^{[n]}(t) = g(g^{[n-1]}(t))$ and $g^{[0]}(t) = t$.) Ozhigov's result is that if L grows at least as fast as M^7 then any quantum algorithm for evaluating the M^{th} iterate takes of order M calls of the unitary operator which evaluates g; of course the classical algorithm requires M calls. Later we will show that our result in fact implies a stronger version of Ozhigov's with L = 2M.

In this paper we show that a quantum computer cannot outperform a classical computer in determining the parity of a function; similar and additional results are obtained in [6] and [7]. Let

$$f(x) = \pm 1 \text{ for } x = 1, \dots N$$
 (1)

Define the parity of f by

$$par(f) = \prod_{x=1}^{N} f(x)$$
(2)

so that the parity of f can be either +1 or -1. The parity of f always depends on the value of f at every point in its domain so classically it requires N function calls to determine the parity. The Grover problem, as described above,

is a special case of the parity problem where additional restrictions have been placed on the function. Although the Grover problem can be solved in time of order \sqrt{N} on a quantum computer, the parity problem has no comparable quantum speedup.

II. PRELIMINARIES

We imagine that the function f whose parity we wish to determine is provided to us in the form of an ordinary computer program, thought of as an oracle. We then use a quantum compiler to convert this to quantum code which gives us the unitary operator

$$U_f|x,+1\rangle = |x,f(x)\rangle$$

$$U_f|x,-1\rangle = |x,-f(x)\rangle.$$
(3)

(Here the second register is a qubit taking the values ± 1 .) Defining

$$|x,s\rangle = \frac{1}{\sqrt{2}}(|x,+1\rangle + |x,-1\rangle)$$

and

$$|x,a\rangle = \frac{1}{\sqrt{2}}(|x,+1\rangle - |x,-1\rangle) , \qquad (4)$$

we have that

$$U_f|x,q\rangle = f(x,q)|x,q\rangle \qquad q = s,a$$
 (5)

where

$$f(x,s) = 1$$
 and $f(x,a) = f(x)$. (6)

Therefore in the $|x,q\rangle$ basis, the quantum operator U_f is multiplication by f(x,q). Suppose that N=2 so that x takes only the values 1 and 2. Then

$$U_f(|1, a\rangle + |2, a\rangle) = f(1)|1, a\rangle + f(2)|2, a\rangle = f(1)(|1, a\rangle + par(f)|2, a\rangle).$$
 (7)

Now the states $|1,a\rangle + |2,a\rangle$ and $|1,a\rangle - |2,a\rangle$ are orthogonal so we see that one application of U_f determines the parity of f although classically two function calls are required. See for example [8]. In section IV this algorithm is generalized for the case of N to determine parity after N/2 applications of U_f .

In writing (3) we ignored the work bits used in calculating f(x). This is because, quite generally, the work bits can be reset to their x independent values [9]. To do this you must first copy f(x) and then run the quantum algorithm for evaluating f(x) backwards thereby resetting the work bits. If this is done then a single application of U_f can be counted as two calls of f.

III. MAIN RESULT

We imagine that we have a quantum algorithm for determining the parity of a function f. The Hilbert space we are working in may be much larger than the 2N-dimensional space spanned by the vectors $|x,q\rangle$ previously described. The algorithm is a sequence of unitary operators which acts on an initial vector $|\psi_0\rangle$ and produces $|\psi_f\rangle$. The Hilbert space is divided into two orthogonal subspaces by a projection operator \mathcal{P} . After producing $|\psi_f\rangle$, we measure \mathcal{P} obtaining either 0, corresponding to parity -1, or 1, corresponding to parity +1. (Note that $\langle \psi_f | \mathcal{P} | \psi_f \rangle$ is the probability of obtaining 1.) We say that the algorithm is successful if there is an $\epsilon > 0$ such that

For par
$$(f) = +1$$
, $\langle \psi_f | \mathcal{P} | \psi_f \rangle \ge \frac{1}{2} + \epsilon$

and

For par
$$(g) = -1$$
, $\langle \psi_g | \mathcal{P} | \psi_g \rangle \leq \frac{1}{2} - \epsilon$. (8)

This is a weak definition of success for an algorithm—we only ask that the probability of correctly identifying the parity of f be greater than $\frac{1}{2}$ no matter what f is. Since we are proving the nonexistence of a successful (short) algorithm, our result is correspondingly strong.

The algorithm is a sequence of unitary operators, some of which are independent of f, and some of which depend on f through the application of a generalization of (5). We need to generalize (5) because we are working in a larger Hilbert space. In this larger Hilbert space there are still subspaces associated with x and q. (In other words, there is a basis of the form $|x,q,w\rangle$ where $x=1,\ldots N$ and q=a,s and $w=1,\ldots W$ for some W, corresponding to the values of the work bits that the algorithm may use.) Accordingly there are projection operators P_x and P_q which obey

$$P_x^2 = P_x \; ; \quad P_x P_y = 0 \text{ for } x \neq y \; ; \quad \sum_{x=1}^N P_x = 1$$

and

$$P_q^2 = P_q \; ; \quad P_s P_a = 0 \; ; \quad \sum_{q=s,a} P_q = 1 \; .$$
 (9)

In terms of these projectors we have

$$U_f = \sum_{x} \sum_{q} f(x, q) P_x P_q \tag{10}$$

where the sum over x is from 1 to N and the q sum is over s and a.

An algorithm which contains k applications of U_f , acting on $|\psi_0\rangle$, produces

$$|\psi_f\rangle = V_k U_f V_{k-1} U_f \dots V_1 U_f |\psi_0\rangle \tag{11}$$

where V_1 through V_k are unitary operators independent of f, but which may involve the work bits. For more extensive discussion, see [10].

We will now use (10) to put $\langle \psi_f | \mathcal{P} | \psi_f \rangle$ in a form where we can see explicitly how it depends on f, allowing us to show that (8) is impossible if k is too small. We have

$$\langle \psi_f | \mathcal{P} | \psi_f \rangle = \sum_{x_1 q_1} \sum_{x_2 q_2} \dots \sum_{x_{2k} q_{2k}} A(x_1, q_1 \dots x_{2k}, q_{2k}) \prod_{i=1}^{2k} f(x_i, q_i)$$
(12)

where

$$A(x_1, q_1 \dots x_{2k}, q_{2k}) = \langle \psi_0 | P_{x_1} P_{q_1} V_1^{\dagger} \dots V_k^{\dagger} \mathcal{P} V_k \dots V_1 P_{x_{2k}} P_{q_{2k}} | \psi_0 \rangle . \tag{13}$$

Note that A does not depend on f.

There are 2^N different possible f's of the form given by (1). We now sum over all these functions and compute

$$\sum_{f} \langle \psi_f | \mathcal{P} | \psi_f \rangle \operatorname{par}(f) = \sum_{f} \sum_{x_1 q_1} \dots \sum_{x_{2k} q_{2k}} A(x_1, q_1 \dots x_{2k}, q_{2k}) \prod_{i=1}^{2k} f(x_i, q_i) \prod_{y=1}^{N} f(y) . \tag{14}$$

Note that

$$\sum_{f} f(z) = 0 \quad \text{for} \quad z = 1, \dots N$$
(15)

because for each function with f(z) = +1 there is a function with f(z) = -1. Similarly if $z_1, z_2 \dots z_n$ are all distinct, we have

$$\sum_{f} f(z_1) f(z_2) \dots f(z_n) = 0 . {16}$$

Return to (14) and consider the sum on f,

$$\sum_{f} \prod_{i=1}^{2k} f(x_i, q_i) \prod_{y=1}^{N} f(y)$$
 (17)

where $x_1, x_2 \dots x_{2k}$ and $q_1, q_2 \dots q_{2k}$ are fixed. For any i with $q_i = s$ we have $f(x_i, s) = 1$. Thus (17) equals

$$\sum_{f} \prod_{\substack{i \text{ with} \\ q_i = a}} f(x_i) \prod_{y=1}^{N} f(y) . \tag{18}$$

Now $f^2(z) = 1$ for any z and any f. By (16), the sum over f in (18) will give 0 unless each term in the second product can be matched to a term in the first product. Since the first product has at most 2k terms and the second product has N terms, we see that if 2k < N then the sum over f in (18) is 0 and accordingly,

$$\sum_{f} \langle \psi_f | \mathcal{P} | \psi_f \rangle \operatorname{par}(f) = 0 . \tag{19}$$

This implies that for 2k < N

$$\sum_{f, \text{par}(f)=+1} \langle \psi_f | \mathcal{P} | \psi_f \rangle = \sum_{f, \text{par}(f)=-1} \langle \psi_f | \mathcal{P} | \psi_f \rangle$$
(20)

which means that for k < N/2 condition (8) cannot be fulfilled.

Equation (20) shows that our bound holds even if we further relax the success criterion given in condition (8). In any algorithm with fewer than N/2 applications of U_f , demanding a probability of success greater than or equal to 1/2 for every f forces the probability to be 1/2 for every f.

IV. AN OPTIMAL ALGORITHM

To see that the bound k < N/2 is optimal, we now show how to solve the parity problem with N/2 applications of U_f . Here we assume that N is even. We only need the states $|x,a\rangle$ given in (4) for which

$$U_f|x,a\rangle = f(x)|x,a\rangle. \tag{21}$$

Define

$$V|x,a\rangle = |x+1,a\rangle \qquad x = 1, \dots \frac{N}{2} - 1$$

$$V|\frac{N}{2},a\rangle = |1,a\rangle$$

$$V|x,a\rangle = |x+1,a\rangle \qquad x = \frac{N}{2} + 1, \dots N - 1$$

$$V|N,a\rangle = |\frac{N}{2} + 1,a\rangle$$
(22)

Also let

$$|\psi_0\rangle = \frac{1}{\sqrt{N}} \sum_{x=1}^{N} |x, a\rangle . \tag{23}$$

Now compute $|\psi_f\rangle$ given by (11) with k=N/2 and for the operators independent of f take

$$V_1 = V_2 = \ldots = V_{k-1} = V$$
 and $V_k = 1$.

We then have that

$$|\psi_f\rangle = \frac{1}{\sqrt{N}}f(1)f(2)\dots f(\frac{N}{2})\sum_{x=1}^{N/2}|x,a\rangle + \frac{1}{\sqrt{N}}f(\frac{N}{2}+1)f(\frac{N}{2}+2)\dots f(N)\sum_{x=\frac{N}{2}+1}^{N}|x,a\rangle.$$
 (24)

Therefore if par (f) = +1, the state $|\psi_f\rangle$ is proportional to $|\psi_0\rangle$ whereas if par (f) = -1, then $|\psi_f\rangle$ is orthogonal to $|\psi_0\rangle$. For the parity projection operator we take $\mathcal{P} = |\psi_0\rangle\langle\psi_0|$ and we see that the algorithm determines the correct parity all the time. Similarly we can show that if N is odd, then with k = (N+1)/2 applications of U_f we can determine the parity of f, but this time we need the states $|x,s\rangle$ as well as $|x,a\rangle$.

V. PARITY AS ITERATED FUNCTION EVALUATION

Here we are interested in evaluating the N^{th} iterate of a function which maps a set of size 2N to itself. We show that it is impossible for a quantum computer to solve this problem with fewer than N/2 applications of the unitary operator corresponding to the function. As noted above, this is a considerable strengthening of Ozhigov's result.

We assume an algorithm satisfying the above conditions exists and we obtain a contradiction. Let the set of 2N elements be $\{(x,r)\}$ where $x=1,\ldots N$ and $r=\pm 1$. For any f of the form (1) define

$$g(x,r) = (x+1, rf(x))$$
 (25)

where we interpret N+1 as 1. Note that

$$q^{[N]}(1,1) = (1, par(f)). \tag{26}$$

Thus an algorithm which computes the N^{th} iterate of g with fewer than N/2 applications of the corresponding unitary operator would in fact solve the parity problem impossibly fast.

VI. CONCLUSION

Grover's result raised the possibility that any problem involving a function with N inputs could be solved quantum mechanically with only \sqrt{N} applications of the corresponding operator. We have shown that this is not the case. For the parity problem, N/2 applications of the quantum operator are required.

Acknowledgment

Three of us are grateful to the fourth.

- This work was supported in part by The Department of Energy under cooperative agreement DE-FC02-94ER40818 and by the National Science Foundation under grant NSF 95–03322 CCR.
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